

On the IGMF strength and filling factor

Sergey Ostapchenko (NTNU)

3rd Fermi Symposium, May 9-12 2011, Rome

in collaboration with K. Dolag, M. Kachelriess, and R. Tomas

arXiv:1009.1782

Intergalactic Magnetic Field

- magnetic fields in galaxies/galaxy clusters:
from amplification of (much weaker) seed fields

Intergalactic Magnetic Field

- magnetic fields in galaxies/galaxy clusters:
from amplification of (much weaker) seed fields
- **initial seed fields**
 - created in the early universe (e.g. during phase transitions)
 - or created by early starburst galaxies/AGNs
 - or created by recent AGNs (and concentrated in filaments)

Intergalactic Magnetic Field

- magnetic fields in galaxies/galaxy clusters:
from amplification of (much weaker) seed fields
- initial seed fields
 - created in the early universe (e.g. during phase transitions)
 - or created by early starburst galaxies/AGNs
 - or created by recent AGNs (and concentrated in filaments)
- **only weak upper limits exist from Faraday RM**

Intergalactic Magnetic Field

- magnetic fields in galaxies/galaxy clusters:
from amplification of (much weaker) seed fields
- initial seed fields
 - created in the early universe (e.g. during phase transitions)
 - or created by early starburst galaxies/AGNs
 - or created by recent AGNs (and concentrated in filaments)
- only weak upper limits exist from Faraday RM
- **alternative way – studies of TeV γ -rays:**
 - delayed “echoes” of γ -ray flares (Plaga 1995)
 - formation of γ -ray “halos” around point-like sources (Aharonian, Coppi & Völk 1994)

Intergalactic Magnetic Field

- magnetic fields in galaxies/galaxy clusters:
from amplification of (much weaker) seed fields
- initial seed fields
 - created in the early universe (e.g. during phase transitions)
 - or created by early starburst galaxies/AGNs
 - or created by recent AGNs (and concentrated in filaments)
- only weak upper limits exist from Faraday RM
- alternative way – studies of TeV γ -rays:
 - delayed “echoes” of γ -ray flares (Plaga 1995)
 - formation of γ -ray “halos” around point-like sources (Aharonian, Coppi & Völk 1994)
- new: **limits on IGMF from non-observation of GeV γ -rays from TeV blazars** (Neronov & Vovk 2010; Tavecchio et al. 2010)

Limits on IGMF from GeV-TeV observations of blazars

- **limits on the strength of IGMF from GeV-silent TeV blazars** (Neronov & Vovk 2010; Tavecchio et al. 2010)
 - TeV γ -rays pair-produce on EBL photons \Rightarrow e/m cascades
 - results in significant fluxes of secondary γ in the GeV range
 - stationary sources at large z , with hard TeV spectrum and low intrinsic GeV spectrum (e.g. 1ES 0229+200):
cascade γ -s should be observable with Fermi-LAT

Limits on IGMF from GeV-TeV observations of blazars

- limits on the strength of IGMF from GeV-silent TeV blazars (Neronov & Vovk 2010; Tavecchio et al. 2010)
 - TeV γ -rays pair-produce on EBL photons \Rightarrow e/m cascades
 - results in significant fluxes of secondary γ in the GeV range
 - stationary sources at large z , with hard TeV spectrum and low intrinsic GeV spectrum (e.g. 1ES 0229+200): cascade γ -s should be observable with Fermi-LAT
- non-observation of GeV γ -rays \Rightarrow cascade deflections by IGMF
- \Rightarrow allows to get limits on the IGMF strength: $B \gtrsim 10^{-15}$ G

Limits on IGMF from GeV-TeV observations of blazars

- limits on the strength of IGMF from GeV-silent TeV blazars (Neronov & Vovk 2010; Tavecchio et al. 2010)
 - TeV γ -rays pair-produce on EBL photons \Rightarrow e/m cascades
 - results in significant fluxes of secondary γ in the GeV range
 - stationary sources at large z , with hard TeV spectrum and low intrinsic GeV spectrum (e.g. 1ES 0229+200): cascade γ -s should be observable with Fermi-LAT
- non-observation of GeV γ -rays \Rightarrow cascade deflections by IGMF
- \Rightarrow allows to get limits on the IGMF strength: $B \gtrsim 10^{-15}$ G
- open questions:
 - potential source variability?
 - impact of IGMF spacial structure

MC approach: EIMag code

- **e/m cascade** on background photons:
 - pair production: $\gamma\gamma_b \rightarrow e^+e^-$
 - ICS: $e^\pm\gamma_b \rightarrow e^\pm\gamma$
 - synchrotron energy loss for e^\pm

MC approach: EIMag code

- e/m cascade on background photons:
 - pair production: $\gamma\gamma_b \rightarrow e^+e^-$
 - ICS: $e^\pm\gamma_b \rightarrow e^\pm\gamma$
 - synchrotron energy loss for e^\pm
- (1+1)-dimensional treatment
 - production angles neglected
 - deflection in IGMF accounted for (small angle approximation)

MC approach: EIMag code

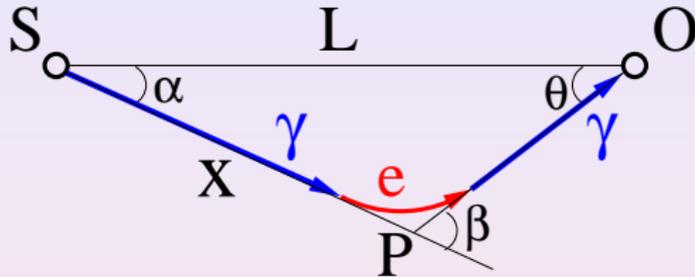
- e/m cascade on background photons:
 - pair production: $\gamma\gamma_b \rightarrow e^+e^-$
 - ICS: $e^\pm\gamma_b \rightarrow e^\pm\gamma$
 - synchrotron energy loss for e^\pm
- (1+1)-dimensional treatment
 - production angles neglected
 - deflection in IGMF accounted for (small angle approximation)
- **weighted sampling applied**
 - produced particle kept with probability $z_E^{\alpha_w}$ ($0 < \alpha_w \leq 1$)
 - each particle is “representative” – has a weight factor w
 - produced particle weight: $w_d = w_p/z_E^{\alpha_w}$

MC approach: EIMag code

- e/m cascade on background photons:
 - pair production: $\gamma\gamma_b \rightarrow e^+e^-$
 - ICS: $e^\pm\gamma_b \rightarrow e^\pm\gamma$
 - synchrotron energy loss for e^\pm
- (1+1)-dimensional treatment
 - production angles neglected
 - deflection in IGMF accounted for (small angle approximation)
- weighted sampling applied
 - produced particle kept with probability $z_E^{\alpha_w}$ ($0 < \alpha_w \leq 1$)
 - each particle is “representative” – has a weight factor w
 - produced particle weight: $w_d = w_p/z_E^{\alpha_w}$
- **highly efficient:** $\sim 10^3$ cascades/s over cosmological distances

Angular deflections in magnetic fields

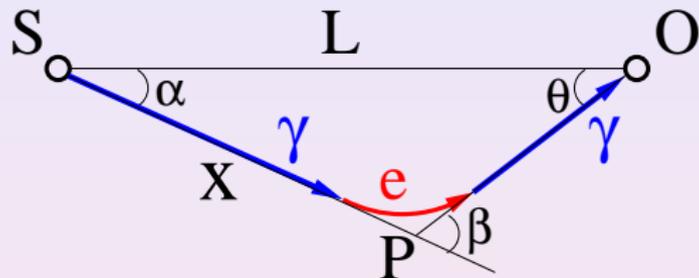
- simple for 2-step process: $\gamma \rightarrow e^\pm \rightarrow \gamma$



- $\beta \equiv \vartheta_{\text{defl}}$ – defl. angle
- $\vartheta \equiv \vartheta_{\text{obs}}$ – obs. angle
- $\beta = \alpha + \vartheta$

Angular deflections in magnetic fields

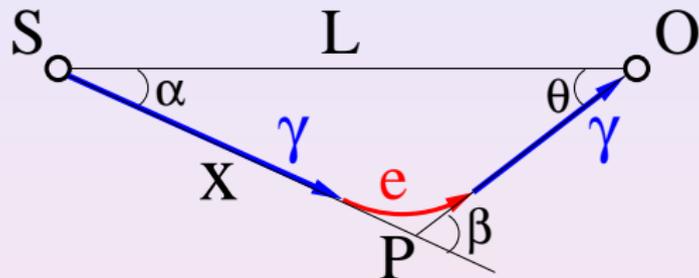
- simple for 2-step process: $\gamma \rightarrow e^\pm \rightarrow \gamma$



- $\beta \equiv \vartheta_{\text{defl}}$ – defl. angle
- $\vartheta \equiv \vartheta_{\text{obs}}$ – obs. angle
- $\beta = \alpha + \vartheta$
- $\Rightarrow \vartheta_{\text{obs}} \simeq \vartheta_{\text{defl}} x/L$

Angular deflections in magnetic fields

- simple for 2-step process: $\gamma \rightarrow e^\pm \rightarrow \gamma$

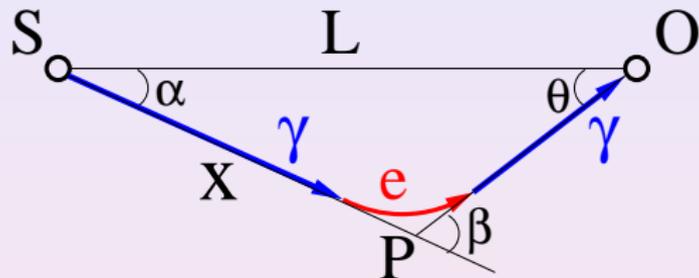


- $\beta \equiv \vartheta_{\text{defl}}$ – defl. angle
- $\vartheta \equiv \vartheta_{\text{obs}}$ – obs. angle
- $\beta = \alpha + \vartheta$
- $\Rightarrow \vartheta_{\text{obs}} \simeq \vartheta_{\text{defl}} x/L$

- NB: though $\langle x \rangle = l_{\gamma\gamma_b}$ (m.f.p.), **fluctuations are very important**
 - γ produced close to the source \Rightarrow smaller ϑ_{obs}

Angular deflections in magnetic fields

- simple for 2-step process: $\gamma \rightarrow e^\pm \rightarrow \gamma$

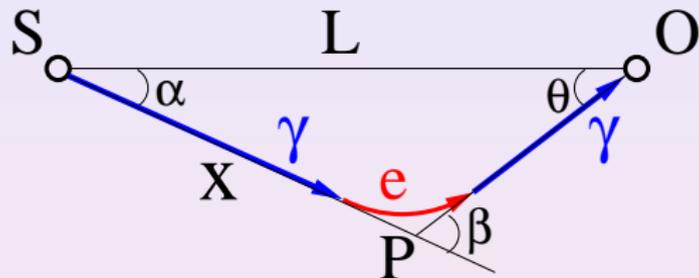


- $\beta \equiv \vartheta_{\text{defl}}$ – defl. angle
- $\vartheta \equiv \vartheta_{\text{obs}}$ – obs. angle
- $\beta = \alpha + \vartheta$
- $\Rightarrow \vartheta_{\text{obs}} \simeq \vartheta_{\text{defl}} x/L$

- NB: though $\langle x \rangle = l_{\gamma\gamma_b}$ (m.f.p.), fluctuations are very important
 - γ produced close to the source \Rightarrow **smaller** ϑ_{obs}

Angular deflections in magnetic fields

- simple for 2-step process: $\gamma \rightarrow e^\pm \rightarrow \gamma$

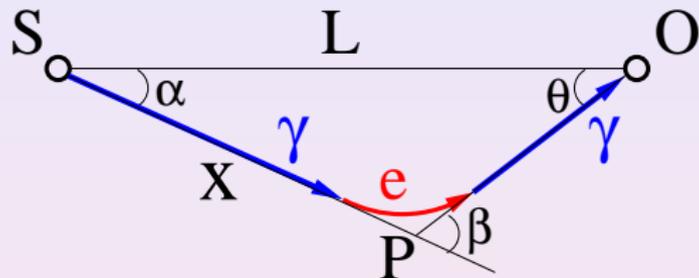


- $\beta \equiv \vartheta_{\text{defl}}$ – defl. angle
- $\vartheta \equiv \vartheta_{\text{obs}}$ – obs. angle
- $\beta = \alpha + \vartheta$
- $\Rightarrow \vartheta_{\text{obs}} \simeq \vartheta_{\text{defl}} x/L$

- NB: though $\langle x \rangle = l_{\gamma\gamma b}$ (m.f.p.), fluctuations are very important
 - γ produced close to the source \Rightarrow smaller ϑ_{obs}
- **time delay:** $\Delta\tau \simeq 2x/c (1 - x/L) \vartheta_{\text{defl}}^2$
 - same importance of fluctuations of x as for ϑ_{obs}

Angular deflections in magnetic fields

- simple for 2-step process: $\gamma \rightarrow e^\pm \rightarrow \gamma$

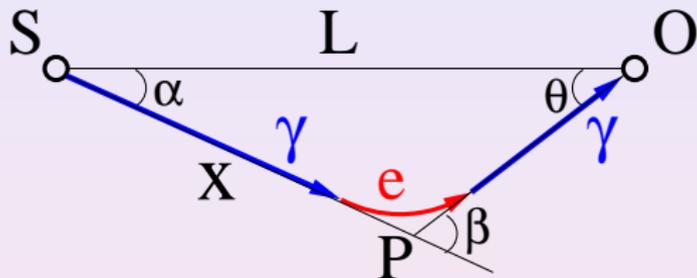


- $\beta \equiv \vartheta_{\text{defl}}$ – defl. angle
- $\vartheta \equiv \vartheta_{\text{obs}}$ – obs. angle
- $\beta = \alpha + \vartheta$
- $\Rightarrow \vartheta_{\text{obs}} \simeq \vartheta_{\text{defl}} x/L$

- NB: though $\langle x \rangle = l_{\gamma\gamma b}$ (m.f.p.), fluctuations are very important
 - γ produced close to the source \Rightarrow smaller ϑ_{obs}
- time delay: $\Delta\tau \simeq 2x/c (1 - x/L) \vartheta_{\text{defl}}^2$
 - **same importance of fluctuations of x** as for ϑ_{obs}

Angular deflections in magnetic fields

- simple for 2-step process: $\gamma \rightarrow e^\pm \rightarrow \gamma$

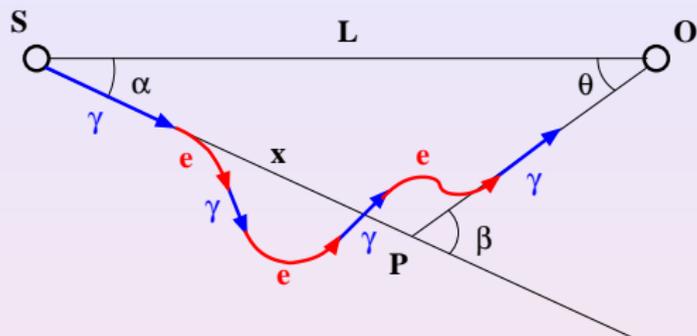


- $\beta \equiv \vartheta_{\text{defl}}$ – defl. angle
- $\vartheta \equiv \vartheta_{\text{obs}}$ – obs. angle
- $\beta = \alpha + \vartheta$
- $\Rightarrow \vartheta_{\text{obs}} \simeq \vartheta_{\text{defl}} x/L$

- NB: though $\langle x \rangle = l_{\gamma\gamma b}$ (m.f.p.), fluctuations are very important
 - γ produced close to the source \Rightarrow smaller ϑ_{obs}
- time delay: $\Delta\tau \simeq 2x/c (1 - x/L) \vartheta_{\text{defl}}^2$
 - same importance of fluctuations of x as for ϑ_{obs}
 - additionally: **fluctuations of Δx_e** ($\langle \Delta x_e \rangle = l_{\text{cool}}$)

Angular deflections in magnetic fields

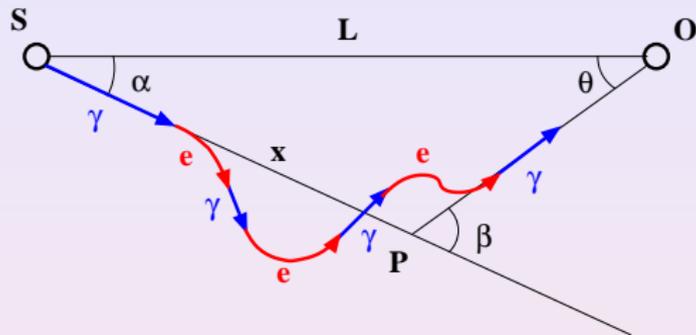
- similarly for multi-step cascades: $\gamma \rightarrow e^\pm \rightarrow \dots \rightarrow e^\pm \rightarrow \gamma$



- $\beta \equiv \vartheta_{\text{defl}}$ – defl. angle
- $\vartheta \equiv \vartheta_{\text{obs}}$ – obs. angle
- $\beta = \alpha + \vartheta$
- $\Rightarrow \vartheta_{\text{obs}} \simeq \vartheta_{\text{defl}} x/L$

Angular deflections in magnetic fields

- similarly for multi-step cascades: $\gamma \rightarrow e^\pm \rightarrow \dots \rightarrow e^\pm \rightarrow \gamma$



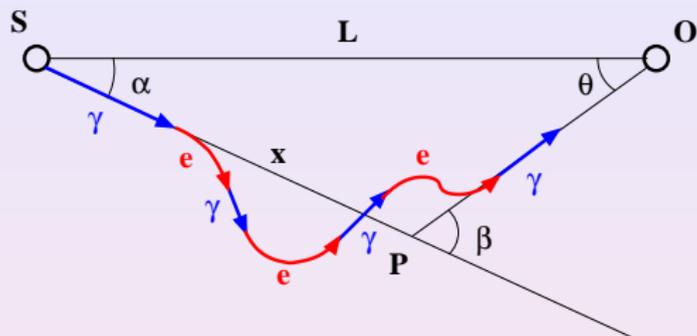
- $\beta \equiv \vartheta_{\text{defl}}$ – defl. angle
- $\vartheta \equiv \vartheta_{\text{obs}}$ – obs. angle
- $\beta = \alpha + \vartheta$
- $\Rightarrow \vartheta_{\text{obs}} \simeq \vartheta_{\text{defl}} x/L$

- deflection angle – within small angle approximation:

- $\vartheta_{\text{defl}} \equiv \beta = \sqrt{\sum_{i=1}^N \beta_{e_i}^2}$ (N – numb. of e^\pm in the cascade branch)
- deflection of the last e^\pm in the cascade – most important (largest x , smallest energy: $\vartheta_{\text{defl}} \sim l_{\text{cool}}/R_L \propto E_e^{-2}$)
- $\beta_{e_i} \sim \Delta x_i$, Δx_i – pass of i -th e^\pm from its creation till γ emission
- if $\Delta x_i \gg L_{\text{coh}} \Rightarrow \beta_{e_i} \sim \sqrt{\Delta x_i}$

Angular deflections in magnetic fields

- similarly for multi-step cascades: $\gamma \rightarrow e^\pm \rightarrow \dots \rightarrow e^\pm \rightarrow \gamma$



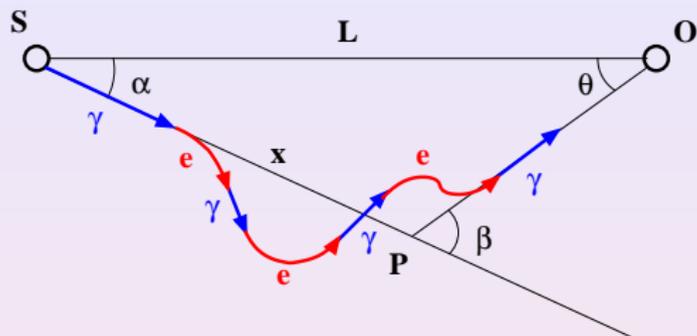
- $\beta \equiv \vartheta_{\text{defl}}$ – defl. angle
- $\vartheta \equiv \vartheta_{\text{obs}}$ – obs. angle
- $\beta = \alpha + \vartheta$
- $\Rightarrow \vartheta_{\text{obs}} \simeq \vartheta_{\text{defl}} x/L$

- deflection angle – within small angle approximation:

- $\vartheta_{\text{defl}} \equiv \beta = \sqrt{\sum_{i=1}^N \beta_{e_i}^2}$ (N – numb. of e^\pm in the cascade branch)
- deflection of the last e^\pm in the cascade** – most important (largest x , smallest energy: $\vartheta_{\text{defl}} \sim l_{\text{cool}}/R_L \propto E_e^{-2}$)
- $\beta_{e_i} \sim \Delta x_i$, Δx_i – pass of i -th e^\pm from its creation till γ emission
- if $\Delta x_i \gg L_{\text{coh}} \Rightarrow \beta_{e_i} \sim \sqrt{\Delta x_i}$

Angular deflections in magnetic fields

- similarly for multi-step cascades: $\gamma \rightarrow e^\pm \rightarrow \dots \rightarrow e^\pm \rightarrow \gamma$



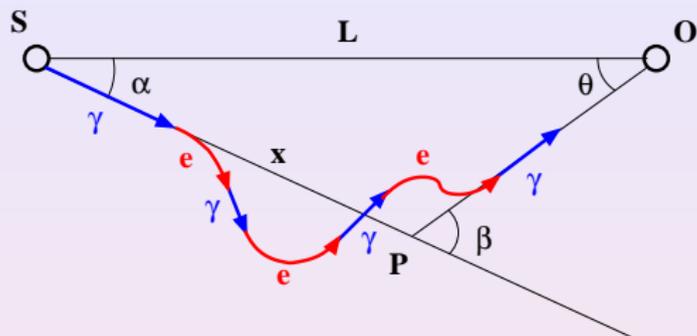
- $\beta \equiv \vartheta_{\text{defl}}$ – defl. angle
- $\vartheta \equiv \vartheta_{\text{obs}}$ – obs. angle
- $\beta = \alpha + \vartheta$
- $\Rightarrow \vartheta_{\text{obs}} \simeq \vartheta_{\text{defl}} x/L$

- deflection angle – within small angle approximation:

- $\vartheta_{\text{defl}} \equiv \beta = \sqrt{\sum_{i=1}^N \beta_{e_i}^2}$ (N – numb. of e^\pm in the cascade branch)
- deflection of the last e^\pm in the cascade – most important (largest x , smallest energy: $\vartheta_{\text{defl}} \sim l_{\text{cool}}/R_L \propto E_e^{-2}$)
- $\beta_{e_i} \sim \Delta x_i$, Δx_i – pass of i -th e^\pm from its creation till γ emission
- if $\Delta x_i \gg L_{\text{coh}} \Rightarrow \beta_{e_i} \sim \sqrt{\Delta x_i}$

Angular deflections in magnetic fields

- similarly for multi-step cascades: $\gamma \rightarrow e^\pm \rightarrow \dots \rightarrow e^\pm \rightarrow \gamma$



- $\beta \equiv \vartheta_{\text{defl}}$ – defl. angle
- $\vartheta \equiv \vartheta_{\text{obs}}$ – obs. angle
- $\beta = \alpha + \vartheta$
- $\Rightarrow \vartheta_{\text{obs}} \simeq \vartheta_{\text{defl}} x/L$

- deflection angle – within small angle approximation:

- $\vartheta_{\text{defl}} \equiv \beta = \sqrt{\sum_{i=1}^N \beta_{e_i}^2}$ (N – numb. of e^\pm in the cascade branch)
- deflection of the last e^\pm in the cascade – most important (largest x , smallest energy: $\vartheta_{\text{defl}} \sim l_{\text{cool}}/R_L \propto E_e^{-2}$)
- $\beta_{e_i} \sim \Delta x_i$, Δx_i – pass of i -th e^\pm from its creation till γ emission
- if $\Delta x_i \gg L_{\text{coh}} \Rightarrow \beta_{e_i} \sim \sqrt{\Delta x_i}$

- same assumptions on the source as in Tavecchio et al. 2010
 - injection spectrum $\mathcal{F} \propto E^{-2/3}$ with cutoff at $E_{\max} = 20$ TeV
 - low Lorentz factor: $\Gamma = 10 \Rightarrow \Theta_{\text{jet}} = 6^\circ$
 - jet pointing towards the observer

Case of 1ES 0229+200: assumptions

- same assumptions on the source as in Tavecchio et al. 2010
 - injection spectrum $\mathcal{F} \propto E^{-2/3}$ with cutoff at $E_{\max} = 20$ TeV
 - low Lorentz factor: $\Gamma = 10 \Rightarrow \Theta_{\text{jet}} = 6^\circ$
 - jet pointing towards the observer

Case of 1ES 0229+200: assumptions

- same assumptions on the source as in Tavecchio et al. 2010
 - injection spectrum $\mathcal{F} \propto E^{-2/3}$ with cutoff at $E_{\max} = 20$ TeV
 - **low Lorentz factor:** $\Gamma = 10 \Rightarrow \Theta_{\text{jet}} = 6^\circ$
 - jet pointing towards the observer

Case of 1ES 0229+200: assumptions

- same assumptions on the source as in Tavecchio et al. 2010
 - injection spectrum $\mathcal{F} \propto E^{-2/3}$ with cutoff at $E_{\max} = 20$ TeV
 - low Lorentz factor: $\Gamma = 10 \Rightarrow \Theta_{\text{jet}} = 6^\circ$
 - jet pointing towards the observer

Case of 1ES 0229+200: assumptions

- same assumptions on the source as in Tavecchio et al. 2010
 - injection spectrum $\mathcal{F} \propto E^{-2/3}$ with cutoff at $E_{\max} = 20$ TeV
 - low Lorentz factor: $\Gamma = 10 \Rightarrow \Theta_{\text{jet}} = 6^\circ$
 - jet pointing towards the observer
- Fermi-LAT upper limits on GeV γ -s from Tavecchio et al. 2010

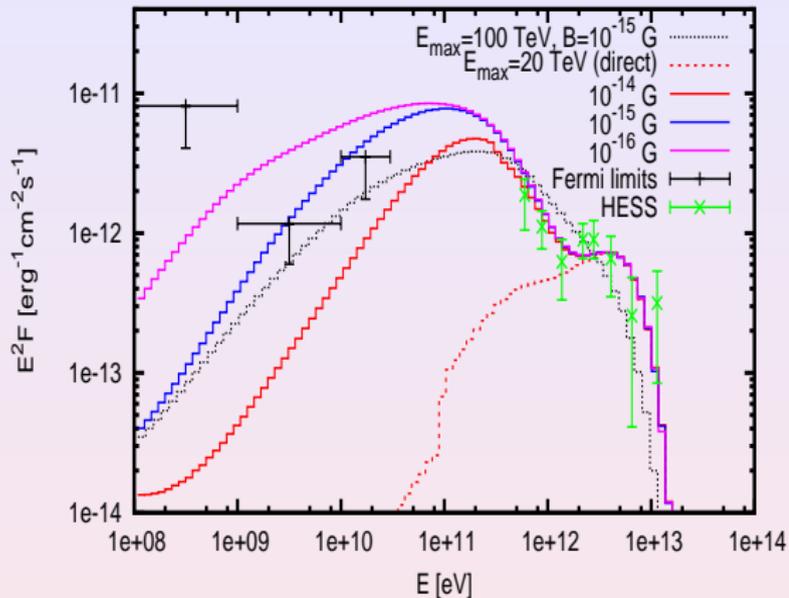
Case of 1ES 0229+200: assumptions

- same assumptions on the source as in Tavecchio et al. 2010
 - injection spectrum $\mathcal{F} \propto E^{-2/3}$ with cutoff at $E_{\max} = 20$ TeV
 - low Lorentz factor: $\Gamma = 10 \Rightarrow \Theta_{\text{jet}} = 6^\circ$
 - jet pointing towards the observer
- Fermi-LAT upper limits on GeV γ -s from Tavecchio et al. 2010
- **account for γ -rays within the PSF of the Fermi-LAT**
($\vartheta_{95} \simeq 1.68^\circ (E/\text{GeV})^{-0.77} + 0.2^\circ \exp(-10 \text{ GeV}/E)$)

Case of 1ES 0229+200: assumptions

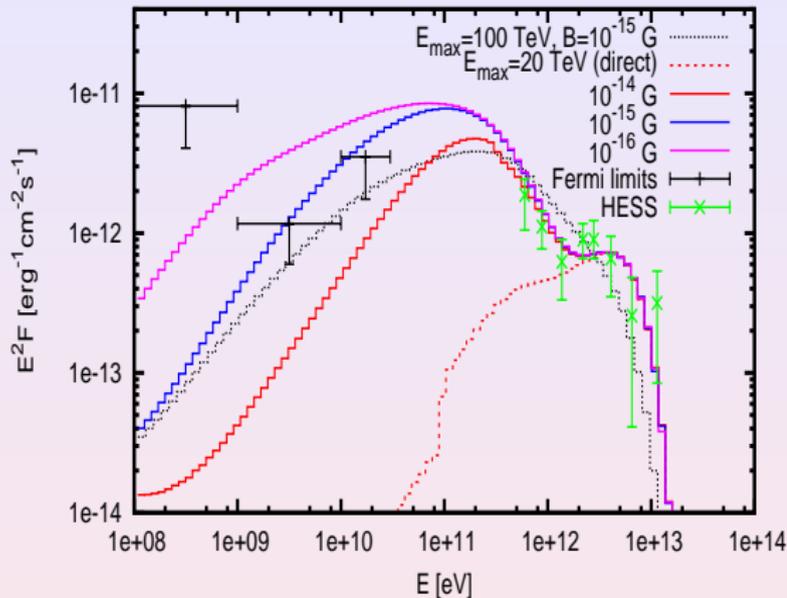
- same assumptions on the source as in Tavecchio et al. 2010
 - injection spectrum $\mathcal{F} \propto E^{-2/3}$ with cutoff at $E_{\max} = 20$ TeV
 - low Lorentz factor: $\Gamma = 10 \Rightarrow \Theta_{\text{jet}} = 6^\circ$
 - jet pointing towards the observer
- Fermi-LAT upper limits on GeV γ -s from Tavecchio et al. 2010
- account for γ -rays within the PSF of the Fermi-LAT
($\vartheta_{95} \simeq 1.68^\circ (E/\text{GeV})^{-0.77} + 0.2^\circ \exp(-10 \text{ GeV}/E)$)
- EBL “best-fit” model from Kneiske & Dole 2010

Case of 1ES 0229+200: γ -ray fluxes



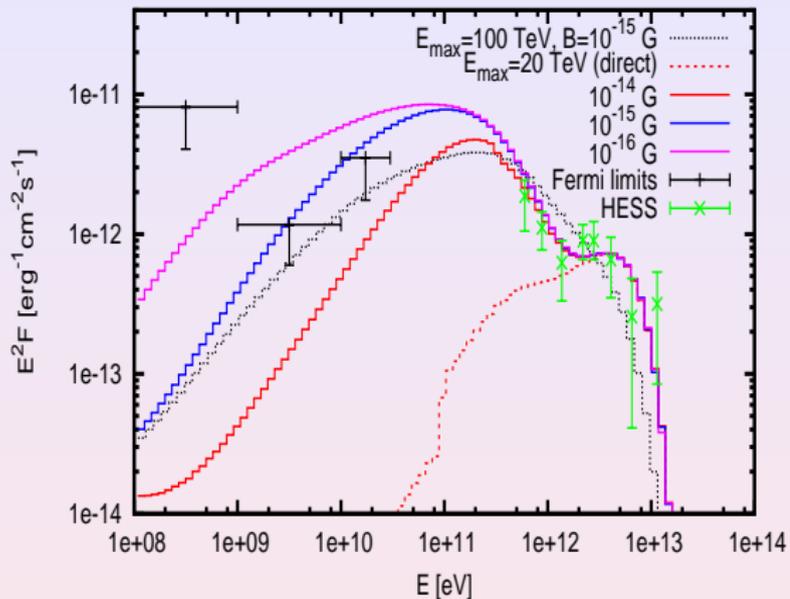
- fluxes normalized to HESS data

Case of 1ES 0229+200: γ -ray fluxes



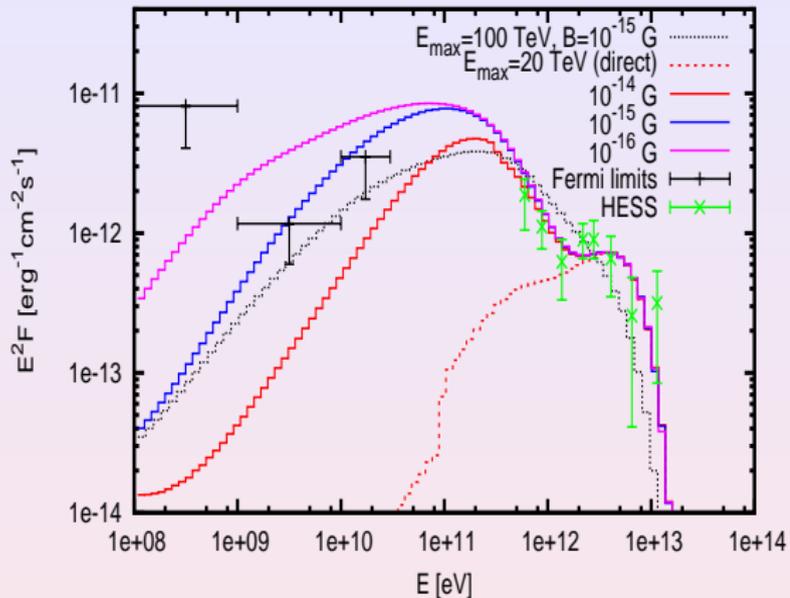
- fluxes normalized to HESS data
- $E_{\text{max}} = 20 \text{ TeV}$:
above Fermi limits
for $B \lesssim 10^{-15} \text{ G}$

Case of 1ES 0229+200: γ -ray fluxes



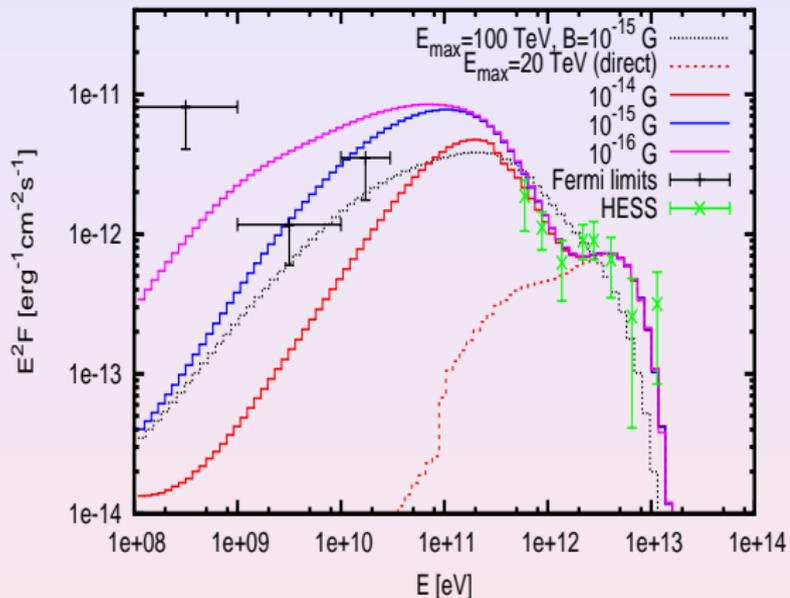
- fluxes normalized to HESS data
- $E_{\max} = 20$ TeV: above Fermi limits for $B \lesssim 10^{-15}$ G
- weaker limits for $E_{\max} = 100$ TeV: 'plato'-like spectra
- $B_{\text{IGMF}} \gtrsim 10^{-15}$ G

Case of 1ES 0229+200: γ -ray fluxes



- fluxes normalized to HESS data
- $E_{\max} = 20$ TeV: above Fermi limits for $B \lesssim 10^{-15}$ G
- weaker limits for $E_{\max} = 100$ TeV: 'plato'-like spectra
- $B_{\text{IGMF}} \gtrsim 10^{-15}$ G

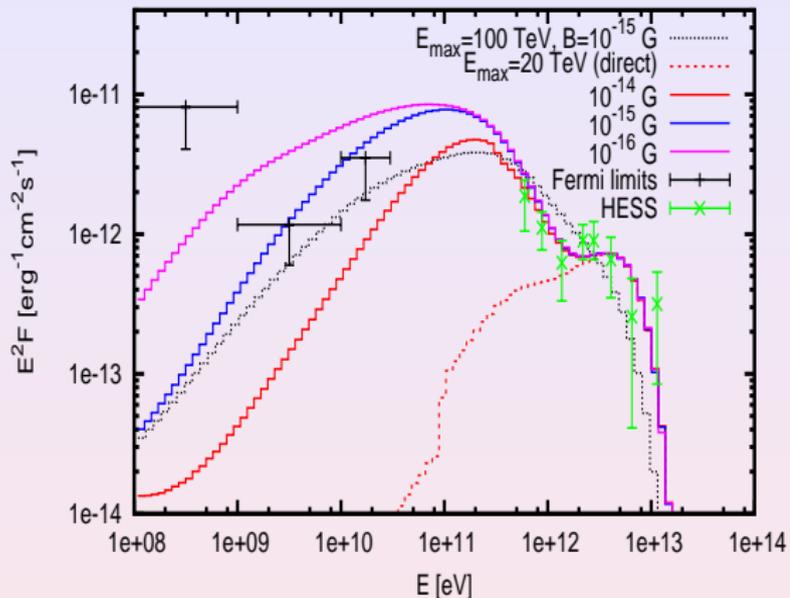
Case of 1ES 0229+200: γ -ray fluxes



- fluxes normalized to HESS data
- $E_{\text{max}} = 20 \text{ TeV}$: above Fermi limits for $B \lesssim 10^{-15} \text{ G}$
- weaker limits for $E_{\text{max}} = 100 \text{ TeV}$: 'plato'-like spectra
- $B_{\text{IGMF}} \gtrsim 10^{-15} \text{ G}$

• results – consistent with Tavecchio et al. 2010

Case of 1ES 0229+200: γ -ray fluxes



- fluxes normalized to HESS data
- $E_{\max} = 20$ TeV: above Fermi limits for $B \lesssim 10^{-15}$ G
- weaker limits for $E_{\max} = 100$ TeV: 'plato'-like spectra
- $B_{\text{IGMF}} \gtrsim 10^{-15}$ G

- results – consistent with Tavecchio et al. 2010

- **different spectral shape**

- e.g. spectral 'shoulder' in the TeV range for $E_{\max} = 20$ TeV

Case of structured magnetic field

- what if the field is concentrated in filaments
while being absent/very weak in voids?

Case of structured magnetic field

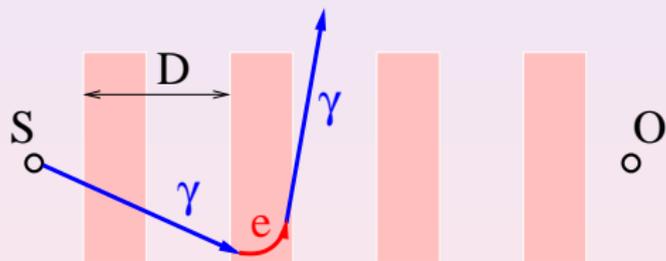
- what if the field is concentrated in filaments while being absent/very weak in voids?
- check with “top-hat” profile, with $D = 10$ Mpc between peaks

Case of structured magnetic field

- what if the field is concentrated in filaments while being absent/very weak in voids?
- check with “top-hat” profile, with $D = 10$ Mpc between peaks
- if the field in “filaments” sufficiently strong \Rightarrow 2 possible cases (since $D \ll l_{\gamma\gamma_b}$, $l_{\text{cool}} \ll (1-f)D$)

Case of structured magnetic field

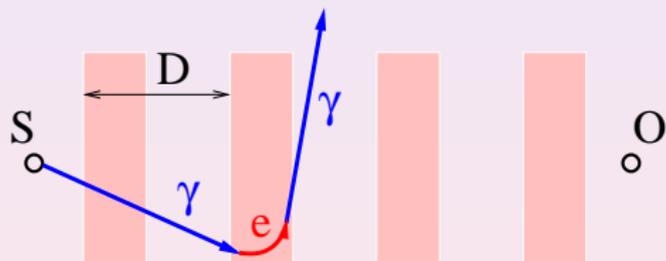
- what if the field is concentrated in filaments while being absent/very weak in voids?
- check with “top-hat” profile, with $D = 10$ Mpc between peaks
- if the field in “filaments” sufficiently strong \Rightarrow 2 possible cases (since $D \ll l_{\gamma\gamma_b}$, $l_{\text{cool}} \ll (1-f)D$)



- with probability f , e^\pm is inside a “filament”
- final γ deflected away

Case of structured magnetic field

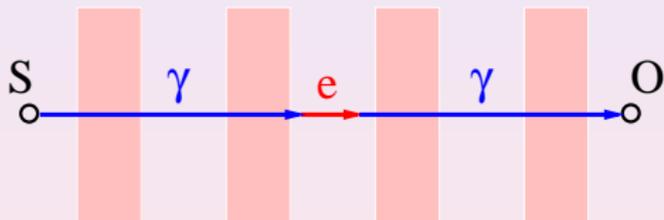
- what if the field is concentrated in filaments while being absent/very weak in voids?
- check with “top-hat” profile, with $D = 10$ Mpc between peaks
- if the field in “filaments” sufficiently strong \Rightarrow 2 possible cases (since $D \ll l_{\gamma\gamma_b}$, $l_{\text{cool}} \ll (1-f)D$)



- with probability f , e^\pm is inside a “filament”
- final γ deflected away

Case of structured magnetic field

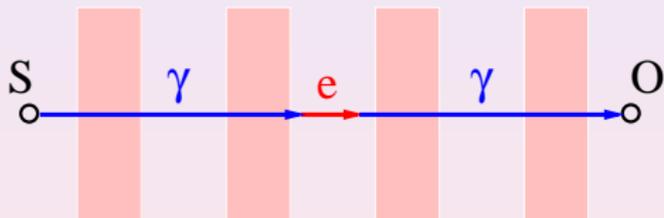
- what if the field is concentrated in filaments while being absent/very weak in voids?
- check with “top-hat” profile, with $D = 10$ Mpc between peaks
- if the field in “filaments” sufficiently strong \Rightarrow 2 possible cases (since $D \ll l_{\gamma b}, l_{\text{cool}} \ll (1-f)D$)



- with probability $(1-f)$, e^\pm produced in a “void”
- \Rightarrow final γ goes straight

Case of structured magnetic field

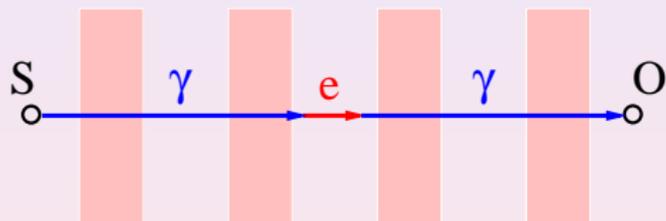
- what if the field is concentrated in filaments while being absent/very weak in voids?
- check with “top-hat” profile, with $D = 10$ Mpc between peaks
- if the field in “filaments” sufficiently strong \Rightarrow 2 possible cases (since $D \ll l_{\gamma\gamma_b}$, $l_{\text{cool}} \ll (1-f)D$)



- with probability $(1-f)$, e^\pm produced in a “void”
- \Rightarrow final γ goes straight

Case of structured magnetic field

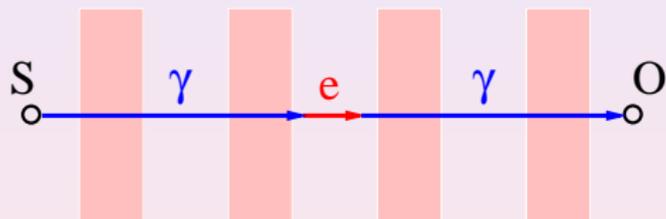
- what if the field is concentrated in filaments while being absent/very weak in voids?
- check with “top-hat” profile, with $D = 10$ Mpc between peaks
- if the field in “filaments” sufficiently strong \Rightarrow 2 possible cases (since $D \ll l_{\gamma b}, l_{\text{cool}} \ll (1-f)D$)



- with probability $(1-f)$, e^\pm produced in a “void”
- \Rightarrow final γ goes straight
- \Rightarrow observed flux = $(1-f) \times \text{flux}(B=0)$
- **multi-step cascade:** observed flux $\sim (1-f)^N \times \text{flux}(B=0)$ (all N electrons in a cascade branch propagate in voids)

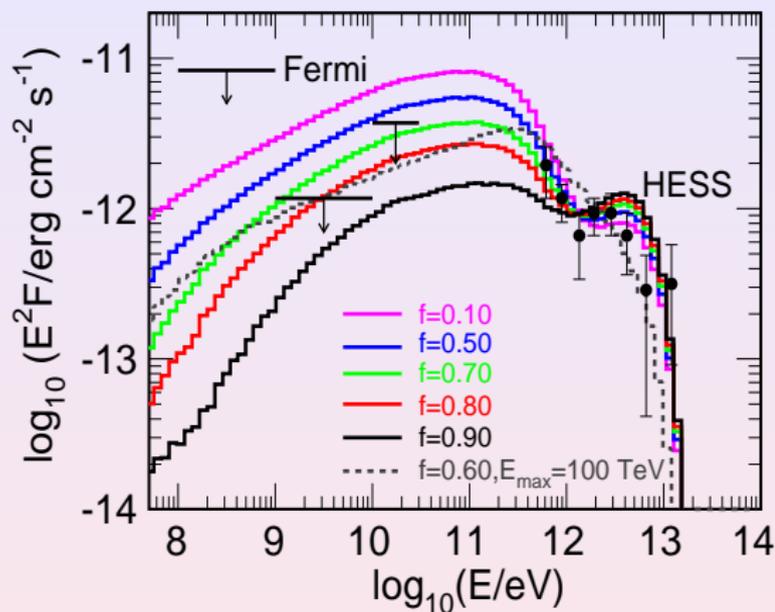
Case of structured magnetic field

- what if the field is concentrated in filaments while being absent/very weak in voids?
- check with “top-hat” profile, with $D = 10$ Mpc between peaks
- if the field in “filaments” sufficiently strong \Rightarrow 2 possible cases (since $D \ll l_{\gamma\gamma_b}, l_{\text{cool}} \ll (1-f)D$)



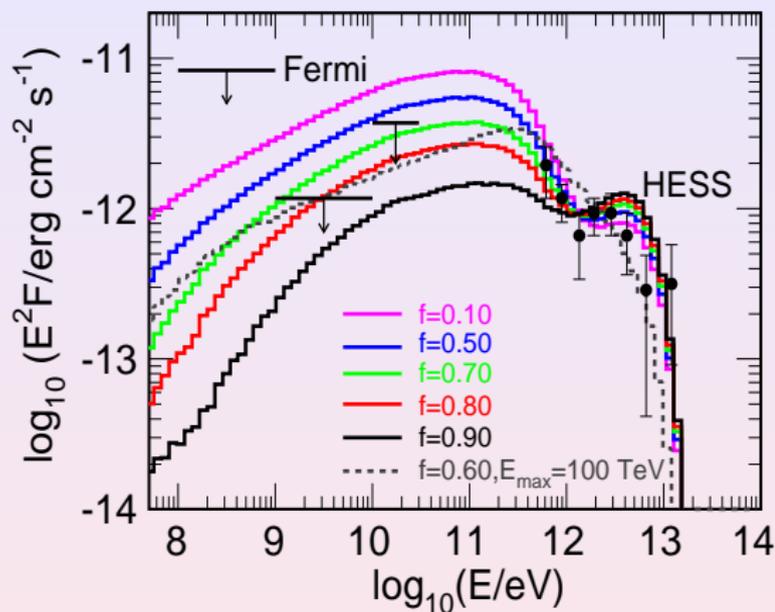
- with probability $(1-f)$, e^\pm produced in a “void”
- \Rightarrow final γ goes straight
- \Rightarrow observed flux = $(1-f) \times \text{flux}(B=0)$
- multi-step cascade: observed flux $\sim (1-f)^N \times \text{flux}(B=0)$ (all N electrons in a cascade branch propagate in voids)
- \Rightarrow lower limit on the “filling factor” – from higher E_{max}

γ -rays from 1ES 0229+200: “filling factor” dependence



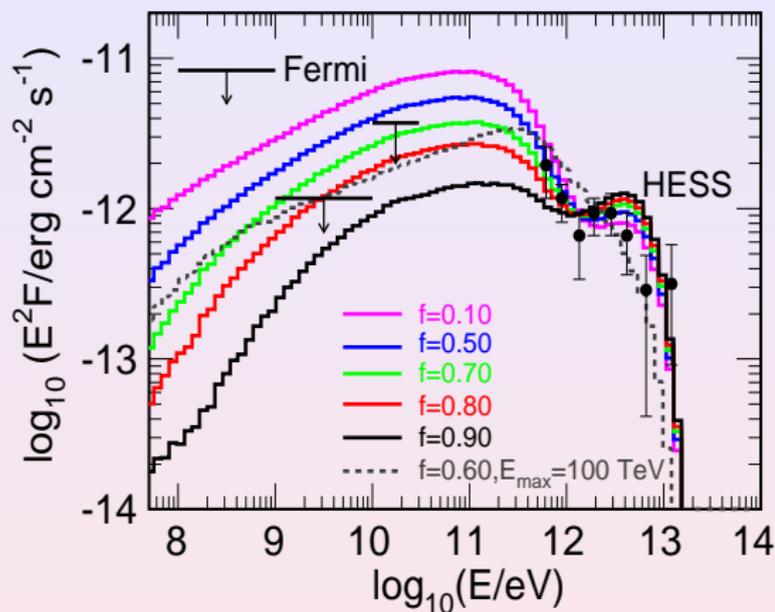
- using $B = 10^{-10} \text{ G}$ in “filaments” ($B = 0$ in “voids”)

γ -rays from 1ES 0229+200: “filling factor” dependence



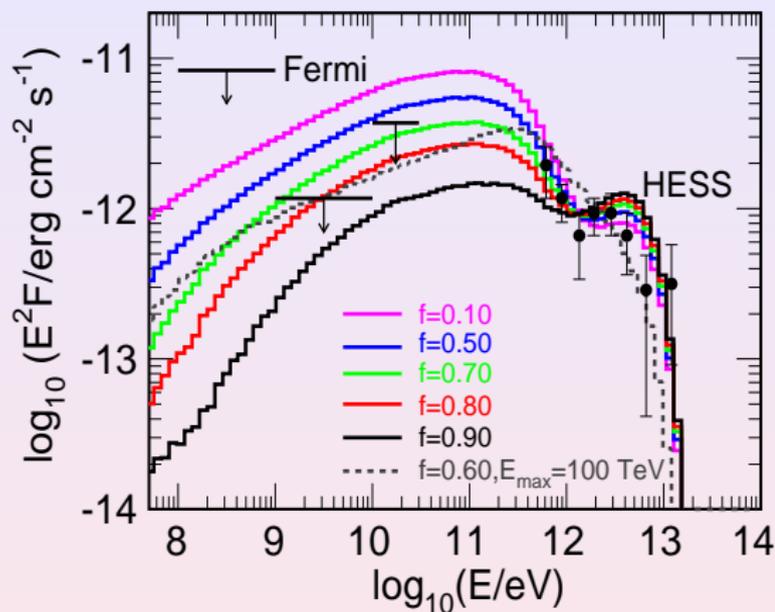
- using $B = 10^{-10} \text{ G}$ in “filaments” ($B = 0$ in “voids”)
- $E_{\text{max}} = 20 \text{ TeV}$:
 $f \gtrsim 0.8$

γ -rays from 1ES 0229+200: “filling factor” dependence



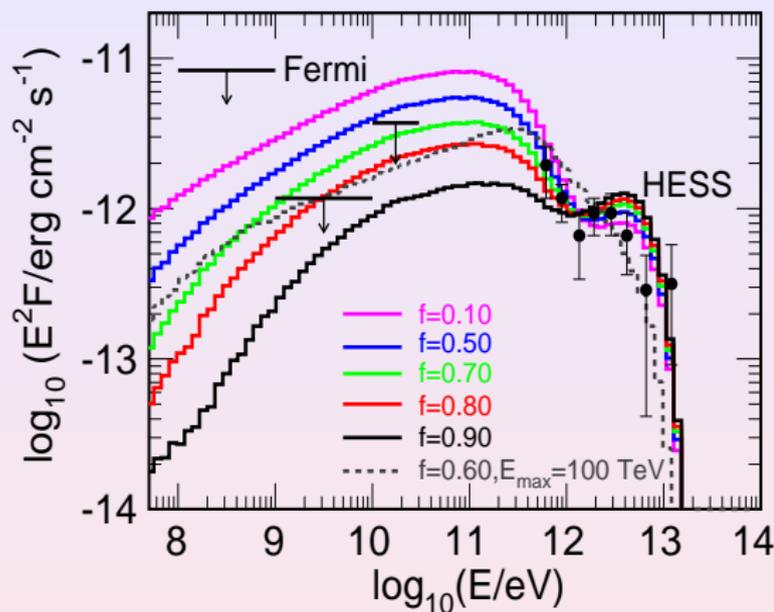
- using $B = 10^{-10}$ G in “filaments” ($B = 0$ in “voids”)
- $E_{\max} = 20$ TeV:
 $f \gtrsim 0.8$
- $E_{\max} = 100$ TeV:
 $f \gtrsim 0.6$

γ -rays from 1ES 0229+200: “filling factor” dependence



- using $B = 10^{-10}$ G in “filaments” ($B = 0$ in “voids”)
- $E_{\max} = 20$ TeV:
 $f \gtrsim 0.8$
- $E_{\max} = 100$ TeV:
 $f \gtrsim 0.6$
- no B -dependence for $B \gtrsim 5 \cdot 10^{-15}$ G

γ -rays from 1ES 0229+200: “filling factor” dependence



- using $B = 10^{-10}$ G in “filaments” ($B = 0$ in “voids”)
- $E_{\max} = 20$ TeV:
 $f \gtrsim 0.8$
- $E_{\max} = 100$ TeV:
 $f \gtrsim 0.6$
- no B -dependence for $B \gtrsim 5 \cdot 10^{-15}$ G

- similar results when using realistic B -profiles from cosmological MHD simulations (Dolag et al., arXiv:1009.1782)

Effect of time delays

- Dermer et al. 2010: time-variability of blazars may significantly weaken the limits on IGMF strength
 - case of 1ES 0229+200 reanalyzed
 - analytic treatment of time delays applied
 - $B_{\text{IGMF}} \gtrsim 10^{-18}$ G obtained

Effect of time delays

- Dermer et al. 2010: time-variability of blazars may significantly weaken the limits on IGMF strength
 - case of 1ES 0229+200 reanalyzed
 - analytic treatment of time delays applied
 - $B_{\text{IGMF}} \gtrsim 10^{-18}$ G obtained

Effect of time delays

- Dermer et al. 2010: time-variability of blazars may significantly weaken the limits on IGMF strength
 - case of 1ES 0229+200 reanalyzed
 - **analytic treatment of time delays applied**
 - $B_{\text{IGMF}} \gtrsim 10^{-18}$ G obtained

Effect of time delays

- Dermer et al. 2010: time-variability of blazars may significantly weaken the limits on IGMF strength
 - case of 1ES 0229+200 reanalyzed
 - analytic treatment of time delays applied
 - $B_{\text{IGMF}} \gtrsim 10^{-18}$ G obtained

Effect of time delays

- Dermer et al. 2010: time-variability of blazars may significantly weaken the limits on IGMF strength
 - case of 1ES 0229+200 reanalyzed
 - analytic treatment of time delays applied
 - $B_{\text{IGMF}} \gtrsim 10^{-18}$ G obtained
- reminding: $\Delta\tau \propto x_\gamma \vartheta_{\text{defl}}^2$
 - \Rightarrow fluctuations of x_γ and Δx_e are important ($\vartheta_{\text{defl}} \propto \Delta x_e$)

Effect of time delays

- Dermer et al. 2010: time-variability of blazars may significantly weaken the limits on IGMF strength
 - case of 1ES 0229+200 reanalyzed
 - analytic treatment of time delays applied
 - $B_{\text{IGMF}} \gtrsim 10^{-18}$ G obtained
- reminding: $\Delta\tau \propto x_\gamma \vartheta_{\text{defl}}^2$
 - \Rightarrow fluctuations of x_γ and Δx_e are important ($\vartheta_{\text{defl}} \propto \Delta x_e$)
- e^\pm emits photons over $l_{\text{cool}} \Rightarrow \langle \Delta\tau \rangle \propto (l_{\text{cool}}/R_L)^2 \propto E_e^{-4} \propto E_\gamma^{-2}$

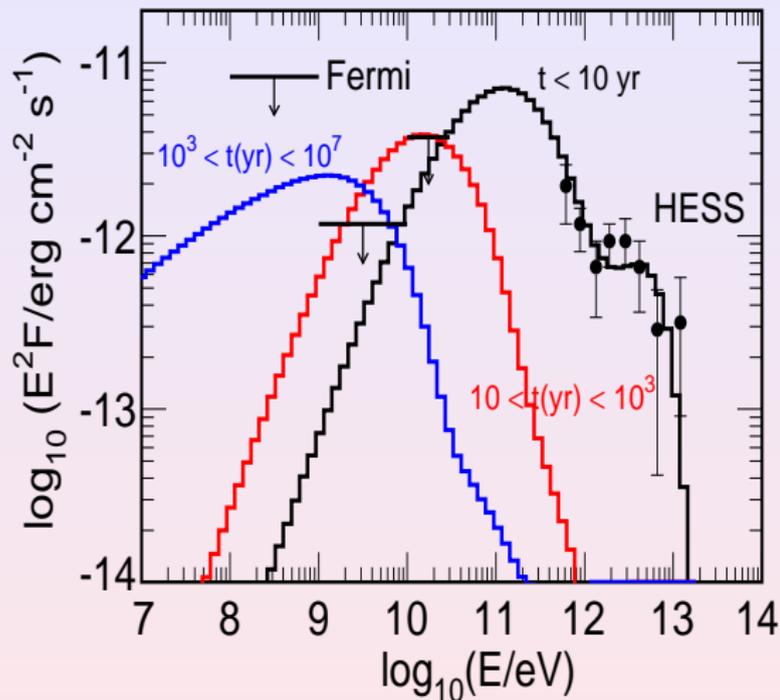
Effect of time delays

- Dermer et al. 2010: time-variability of blazars may significantly weaken the limits on IGMF strength
 - case of 1ES 0229+200 reanalyzed
 - analytic treatment of time delays applied
 - $B_{\text{IGMF}} \gtrsim 10^{-18}$ G obtained
- reminding: $\Delta\tau \propto x_\gamma \vartheta_{\text{defl}}^2$
 - \Rightarrow fluctuations of x_γ and Δx_e are important ($\vartheta_{\text{defl}} \propto \Delta x_e$)
- e^\pm emits photons over $l_{\text{cool}} \Rightarrow \langle \Delta\tau \rangle \propto (l_{\text{cool}}/R_L)^2 \propto E_e^{-4} \propto E_\gamma^{-2}$
- however: **first few photons emitted over $\Delta x_e \sim l_{e\gamma_b} \sim \text{few kpc}$**
 - \Rightarrow contribute to small $\Delta\tau$ ($\sim \langle \Delta\tau \rangle / 100$) for GeV γ -rays

Effect of time delays

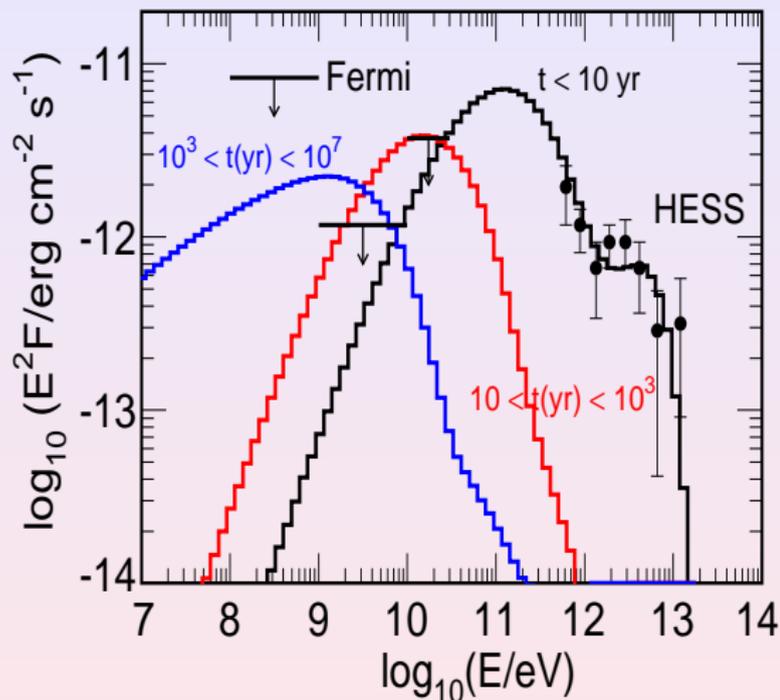
- Dermer et al. 2010: time-variability of blazars may significantly weaken the limits on IGMF strength
 - case of 1ES 0229+200 reanalyzed
 - analytic treatment of time delays applied
 - $B_{\text{IGMF}} \gtrsim 10^{-18}$ G obtained
- reminding: $\Delta\tau \propto x_\gamma \vartheta_{\text{defl}}^2$
 - \Rightarrow fluctuations of x_γ and Δx_e are important ($\vartheta_{\text{defl}} \propto \Delta x_e$)
- e^\pm emits photons over $l_{\text{cool}} \Rightarrow \langle \Delta\tau \rangle \propto (l_{\text{cool}}/R_L)^2 \propto E_e^{-4} \propto E_\gamma^{-2}$
- however: first few photons emitted over $\Delta x_e \sim l_{e\gamma_b} \sim \text{few kpc}$
 - \Rightarrow contribute to small $\Delta\tau$ ($\sim \langle \Delta\tau \rangle / 100$) for GeV γ -rays

γ -rays from 1ES 0229+200: time-dependence



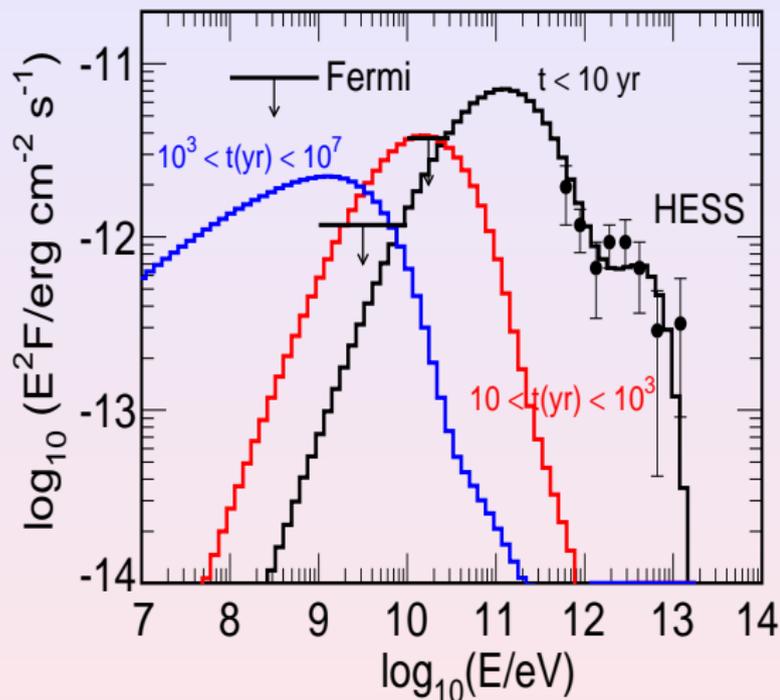
- time-binned spectra
($B = 10^{-17}$ G)

γ -rays from 1ES 0229+200: time-dependence



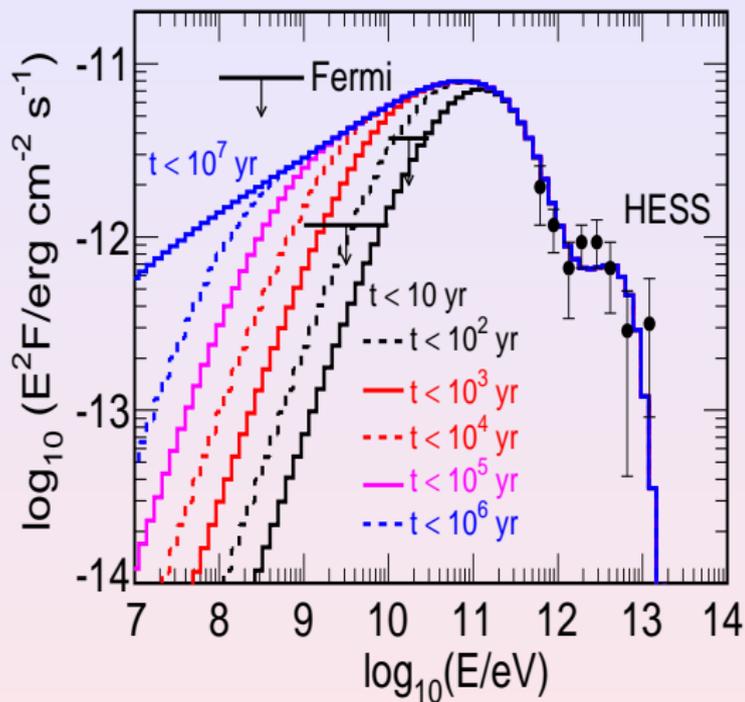
- time-binned spectra ($B = 10^{-17}$ G)
- average time delay: $\langle \Delta\tau \rangle \propto E_\gamma^{-2}$

γ -rays from 1ES 0229+200: time-dependence



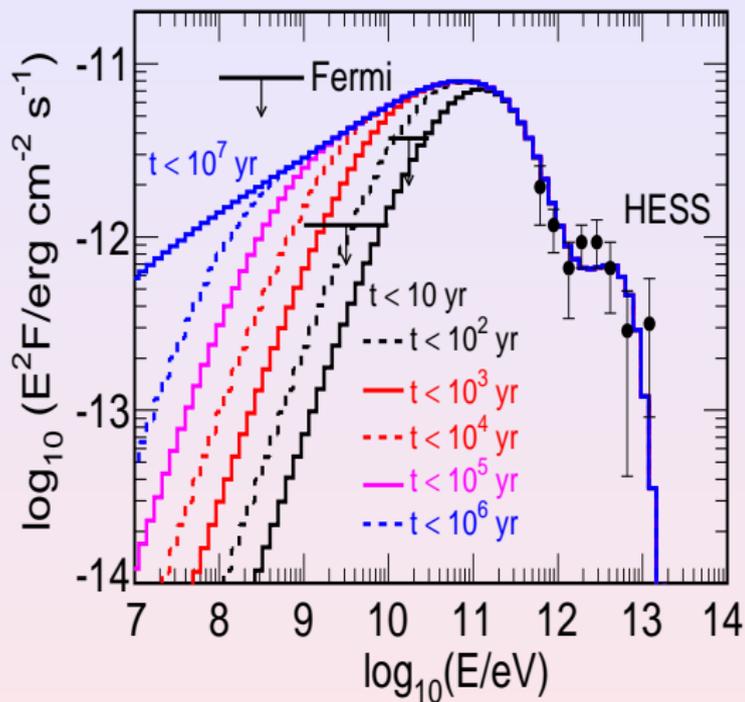
- time-binned spectra ($B = 10^{-17}$ G)
- average time delay: $\langle \Delta\tau \rangle \propto E_\gamma^{-2}$
- distribution shape doesn't scale with E_γ : large width for the $\Delta\tau < 10$ yr bin

γ -rays from 1ES 0229+200: time-dependence



- cumulative spectra
($B = 10^{-17} \text{ G}$)

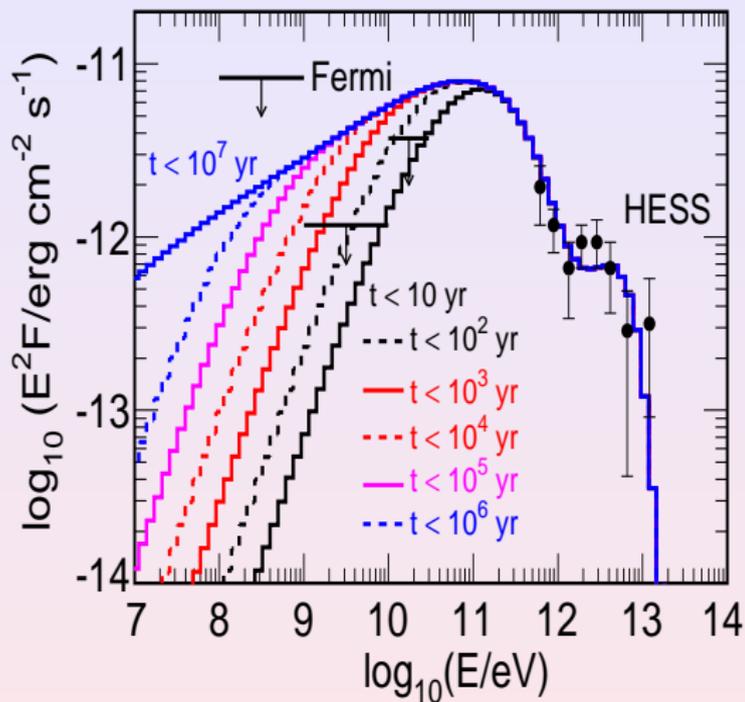
γ -rays from 1ES 0229+200: time-dependence



• cumulative spectra
($B = 10^{-17}$ G)

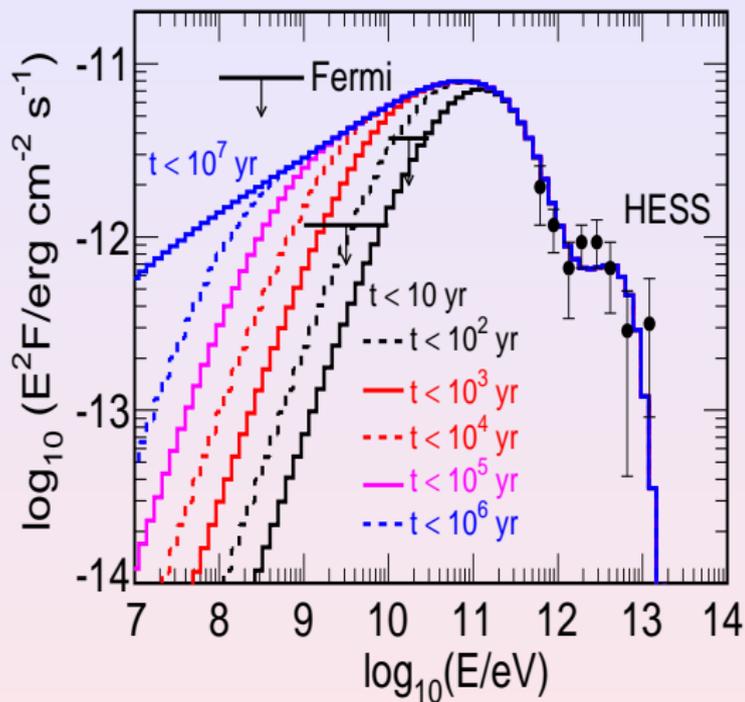
• $\tau_{\text{source}} \gtrsim \text{few yr} \Rightarrow$
 $B \gtrsim 10^{-17}$ G

γ -rays from 1ES 0229+200: time-dependence



- cumulative spectra ($B = 10^{-17}$ G)
- $\tau_{\text{source}} \gtrsim \text{few yr} \Rightarrow B \gtrsim 10^{-17}$ G
- limit scales as $B_{\text{min}} \propto \sqrt{\tau_{\text{source}}}$ (e.g. $B \gtrsim 10^{-15}$ G for $\tau_{\text{source}} \gtrsim \text{few} \times 10^4$ yr)

γ -rays from 1ES 0229+200: time-dependence



- cumulative spectra ($B = 10^{-17}$ G)
- $\tau_{\text{source}} \gtrsim \text{few yr} \Rightarrow B \gtrsim 10^{-17}$ G
- limit scales as $B_{\text{min}} \propto \sqrt{\tau_{\text{source}}}$ (e.g. $B \gtrsim 10^{-15}$ G for $\tau_{\text{source}} \gtrsim \text{few} \times 10^4$ yr)

- similar results obtained by Taylor et al. 2011

limits on IGMF “filling factor”: time-independence

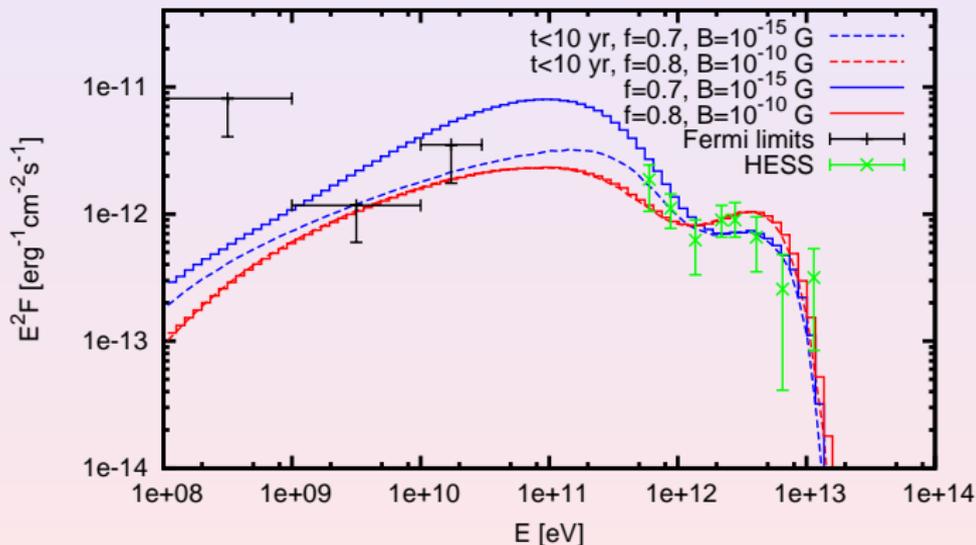
- reminding: “filling factor” related to undeflected γ -ray flux:
observed flux = $(1 - f) \times \text{flux}(B = 0)$

limits on IGMF “filling factor”: time-independence

- reminding: “filling factor” related to undeflected γ -ray flux:
observed flux = $(1 - f) \times \text{flux}(B = 0)$
 - \Rightarrow independent on the life time of the source

limits on IGMF “filling factor”: time-independence

- reminding: “filling factor” related to undeflected γ -ray flux:
observed flux = $(1 - f) \times \text{flux}(B = 0)$
 - \Rightarrow independent on the life time of the source



- variability of the source **impacts the limits on the IGMF strength, not on the IGMF spatial distribution**

Summary

- 1 γ -ray fluxes from 1ES 0229+200 calculated using
MC treatment of e/m cascades on background photons

Summary

- 1 γ -ray fluxes from 1ES 0229+200 calculated using MC treatment of e/m cascades on background photons
- 2 **limits on the IGMF strength and “filling factor” obtained**
 - if the source is stable over $\gtrsim \text{few} \times 10^4$ yr:
fields with $B \gtrsim O(10^{-15})$ G fill more than 60% of space
 - weaker limits on the IGMF strength if the source is variable:
e.g. $B \gtrsim O(10^{-16} \div 10^{-17})$ G for $\tau_{\text{source}} \sim \text{few} \times (100 \div 1)$ yr

Summary

- 1 γ -ray fluxes from 1ES 0229+200 calculated using MC treatment of e/m cascades on background photons
- 2 limits on the IGMF strength and “filling factor” obtained
 - if the source is stable over $\gtrsim \text{few} \times 10^4$ yr:
fields with $B \gtrsim O(10^{-15})$ G fill more than 60% of space
 - weaker limits on the IGMF strength if the source is variable:
e.g. $B \gtrsim O(10^{-16} \div 10^{-17})$ G for $\tau_{\text{source}} \sim \text{few} \times (100 \div 1)$ yr

Summary

- 1 γ -ray fluxes from 1ES 0229+200 calculated using MC treatment of e/m cascades on background photons
- 2 limits on the IGMF strength and “filling factor” obtained
 - if the source is stable over $\gtrsim \text{few} \times 10^4$ yr:
fields with $B \gtrsim O(10^{-15})$ G fill more than 60% of space
 - **weaker limits on the IGMF strength if the source is variable:**
e.g. $B \gtrsim O(10^{-16} \div 10^{-17})$ G for $\tau_{\text{source}} \sim \text{few} \times (100 \div 1)$ yr

Summary

- 1 γ -ray fluxes from 1ES 0229+200 calculated using MC treatment of e/m cascades on background photons
- 2 limits on the IGMF strength and “filling factor” obtained
 - if the source is stable over $\gtrsim \text{few} \times 10^4$ yr:
fields with $B \gtrsim O(10^{-15})$ G fill more than 60% of space
 - weaker limits on the IGMF strength if the source is variable:
e.g. $B \gtrsim O(10^{-16} \div 10^{-17})$ G for $\tau_{\text{source}} \sim \text{few} \times (100 \div 1)$ yr
- 3 limits on the IGMF “filling factor” – independent of the source life time

Summary

- 1 γ -ray fluxes from 1ES 0229+200 calculated using MC treatment of e/m cascades on background photons
- 2 limits on the IGMF strength and “filling factor” obtained
 - if the source is stable over $\gtrsim \text{few} \times 10^4$ yr:
fields with $B \gtrsim O(10^{-15})$ G fill more than 60% of space
 - weaker limits on the IGMF strength if the source is variable:
e.g. $B \gtrsim O(10^{-16} \div 10^{-17})$ G for $\tau_{\text{source}} \sim \text{few} \times (100 \div 1)$ yr
- 3 limits on the IGMF “filling factor” – independent of the source life time
- 4 \Rightarrow strong constraints on the origin of magnetic seed fields
 - volume filling process (e.g. primordial) – strongly favored
 - otherwise: very efficient transport process is required

Summary

- 1 γ -ray fluxes from 1ES 0229+200 calculated using MC treatment of e/m cascades on background photons
- 2 limits on the IGMF strength and “filling factor” obtained
 - if the source is stable over $\gtrsim \text{few} \times 10^4$ yr:
fields with $B \gtrsim O(10^{-15})$ G fill more than 60% of space
 - weaker limits on the IGMF strength if the source is variable:
e.g. $B \gtrsim O(10^{-16} \div 10^{-17})$ G for $\tau_{\text{source}} \sim \text{few} \times (100 \div 1)$ yr
- 3 limits on the IGMF “filling factor” – independent of the source life time
- 4 \Rightarrow strong constraints on the origin of magnetic seed fields
 - **volume filling process (e.g. primordial)** – strongly favored
 - otherwise: very efficient transport process is required

Summary

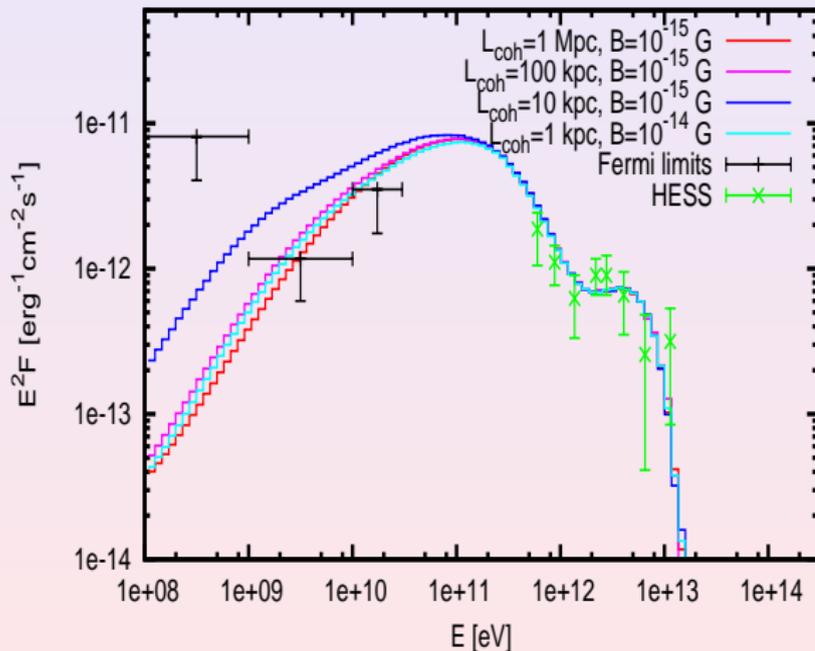
- 1 γ -ray fluxes from 1ES 0229+200 calculated using MC treatment of e/m cascades on background photons
- 2 limits on the IGMF strength and “filling factor” obtained
 - if the source is stable over $\gtrsim \text{few} \times 10^4$ yr:
fields with $B \gtrsim O(10^{-15})$ G fill more than 60% of space
 - weaker limits on the IGMF strength if the source is variable:
e.g. $B \gtrsim O(10^{-16} \div 10^{-17})$ G for $\tau_{\text{source}} \sim \text{few} \times (100 \div 1)$ yr
- 3 limits on the IGMF “filling factor” – independent of the source life time
- 4 \Rightarrow strong constraints on the origin of magnetic seed fields
 - volume filling process (e.g. primordial) – strongly favored
 - otherwise: **very efficient transport process is required**

Backup: coherence length dependence

- for small coherence length of the field **the limit on the IGMF strength improves as $B_{\min} \propto L_{\text{coh}}^{-1/2}$** (Neronov & Semikoz 2009)

Backup: coherence length dependence

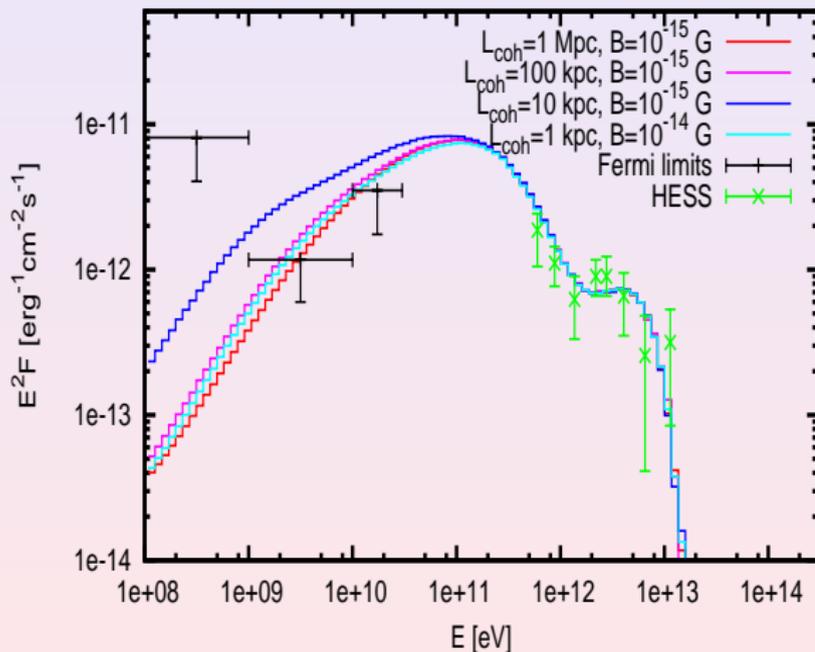
- for small coherence length of the field the limit on the IGMF strength improves as $B_{\min} \propto L_{\text{coh}}^{-1/2}$ (Neronov & Semikoz 2009)



- applies for $L_{\text{coh}} < 100 \text{ kpc}$
- no dependence for $L_{\text{coh}} > 1 \text{ Mpc}$

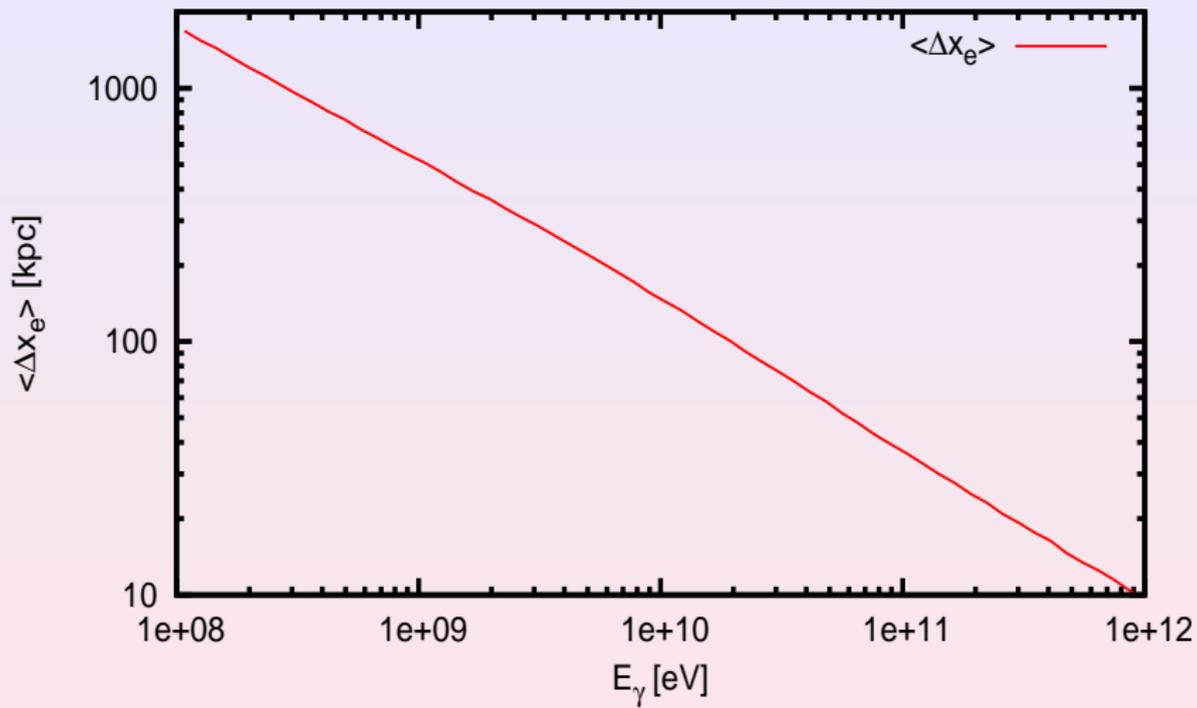
Backup: coherence length dependence

- for small coherence length of the field the limit on the IGMF strength improves as $B_{\min} \propto L_{\text{coh}}^{-1/2}$ (Neronov & Semikoz 2009)



- applies for $L_{\text{coh}} < 100 \text{ kpc}$
- no dependence for $L_{\text{coh}} > 1 \text{ Mpc}$

- mean travel distance $\langle \Delta x_e \rangle$ of a parent e^\pm
is defined by the cooling length



- however, the distribution of Δx_e has pronounced tails towards $\Delta x_e \sim 1$ kpc

